g-2: Status of the Standard Model prediction



Thomas Teubner



- . Introduction
- .. On behalf of Masashi Hayakawa:
 - ullet Status of the 10th order QED contributions to the leptonic g-2
- ... Recent developments in $(g-2)_{\mu}$: Hadronic Vacuum Polarisation contributions
 - 2π channel: Inclusion of Radiative Return data
 - Improvements in the region below 2 GeV
 - pQCD or data for higher energies?
- Full picture. Showdown. $\alpha(M_Z^2)$. Outlook

Thanks to my collaborators Kaoru Hagiwara, Ruofan Liao, Alan Martin and Daisuke Nomura.

$(g-2)_{\mu}$: Contributions

- $a_{\mu} = (g-2)_{\mu}/2 = a_{\mu}^{\text{QED}} + a_{\mu}^{\text{EW}} + a_{\mu}^{\text{had}} + a_{\mu}^{\text{New Physics?}}$
- QED: 4-loop predictions consolidated, 5-loop calculations ongoing, big surprises very unprobable for a_{μ} , error formidably small: $a_{\mu}^{\rm QED}=116584718.08(15)\cdot 10^{-11}$ \checkmark Kinoshita et al.
 - ightarrow However, recent progress at 5-loop level very relevant for a_e

Current status of calculation of 10th-order QED contribution to lepton g-2

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Acknowledgment

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Subject

I overview the current status on the theoretical calculation of the 10th-order (5-loop) QED contribution to the lepton g-2

$$a_{l}(\text{QED}) = \frac{\alpha}{\pi} a_{(1)} + \left(\frac{\alpha}{\pi}\right)^{2} a_{l,(2)}(\text{QED}) + \left(\frac{\alpha}{\pi}\right)^{3} a_{l,(3)}(\text{QED}) + \left(\frac{\alpha}{\pi}\right)^{4} a_{l,(4)}(\text{QED}) + \left(\frac{\alpha}{\pi}\right)^{5} a_{l,(5)}(\text{QED}) + \cdots,$$

and present the *uncertainty* of the most accurate value of α derived using very preliminary result on full $a_{e,(5)}(\text{QED})$ for the electron g-2.

Other works for 10th-order QED contributions

- estimate of logarithmic terms for muon g-2
 - A. L. Kataev, Phys. Rev. D 74, 073011 (2006),
 - P. A. Baikov, K. G. Chetyrkin and C. Sturm, Nucl. Phys. Proc. Suppl. 183, 8 (2008).
- analytic calculation of diagrams with bubbles of one-loop vacuum polarization
 - J. P. Aguilar, D. Greynat and E. De Rafael, Phys. Rev. D 77, 093010 (2008).

Background

Comparison of values of α^{-1} provides a stringent test of our understanding on electromagnetic interaction based on quantum mechanics and field theory:

```
\begin{split} &\alpha^{-1}(\mathrm{Rb06}) = 137\ 035\ 998\ 84\ (91)\ [6.7\mathrm{ppb}]\,,\\ &\alpha^{-1}(\mathrm{Rb11}) = 137\ 035\ 999\ 037\ (91)\ [0.66\mathrm{ppb}]\,,\\ &\alpha^{-1}(\mathrm{a_e08}) = 137\ 035\ 999\ 085\ (12)_8(37)_{10}(33)_{\mathrm{exp}}\ [0.37\mathrm{ppb}]\,, \end{split}
```

where

- $\alpha^{-1}({
 m Rb}11)$ is the latest result by R. Bourchendira, P.Clade, S.Guellati-Khelifa, F.Nez, and F.Biraben, Phys. Rev. Lett. **106**, 080801 (2011) (precise determination of $\hbar/m({
 m Rb})$ using rubidium atom interferometer with Bloch oscillation),
- (continued to the next slide)



- $\alpha^{-1}(a_e08)$ was obtained in T. Aoyama, M. Hayakawa, T. Kinoshita and M. Nio, Phys. Rev. Lett. **99**, 110406 (2007), which corrects the previous 8th-order calculation.
 - (33)_{exp} is the experimental error of $a_e \equiv (g_e 2)/2$, B. C. Odom, D. Hanneke, B. D'Urso and G. Gabrielse, Phys. Rev. Lett. **97**, 030801 (2006),
 - (12)₈ is the theoretical uncertainty associated with the Monte Carlo integration of 8th-order QED correction, $a_{e, (4)}(\text{QED})$,
 - (37)₁₀ is just the guess estimate on 10th-order QED correction, $a_{e,(5)}(\text{QED})$.
 - Seriously, the reduction of the experimental error is scarified due to *lack of our knowledge on a_{e, (5)}* (QED); $(37)_{10} \sim (33)_{\rm exp}$.

Strategy

- Full 10th-order QED calculation needs huge time, which exceeds one's life;
 - About 20 years for 8th order calculation
 ⇒ More than 500 years needed for 10th order calculation...
 - Number of Feynman diagrams = 12,672,
 - Complicated ultra-violet and infrared divergent structures in each diagram
 - ⇒ Difficult to avoid mistakes in renormalization.
- We developed the automation scheme to generate FORTRAN-formatted programs for g - 2 amplitude; Nucl. Phys. B 740, 138 (2006); 796, 184 (2008).
- (continued to the next slide)

- Genericity was discarded for implementation;
 - automation program 1 for the subset of quenched-type Feynman diagrams, whose number amounts to 6,354, about half of the total, 12,672,



- automation program 2 for the subset 2,
- automation program 3 for the subset 3,
- ...

Choice of such a strategy has made 10th-order project in progress quite efficiently.

- Numerical computation has been done using computational resources at RIKEN (RSCC, RICC) over these 6 years (not 500 years), in collaboration with
 - Nio (RIKEN),
 - Asano, Watanabe, Aoyama (Nagoya),
 - Kinoshita (Cornell).

Result

We have a very preliminary result on full 10-th order QED contribution to $g_e - 2$, which gives

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\begin{split} \alpha^{-1}(\mathrm{Rb06}) &= 137\ 035\ 998\ 84\ (91)\ [6.7\mathrm{ppb}]\,,\\ \alpha^{-1}(\mathrm{Rb11}) &= 137\ 035\ 999\ 037\ (91)\ [0.66\mathrm{ppb}]\,,\\ \alpha^{-1}(\mathrm{a_e}08) &= 137\ 035\ 999\ 085\ (12)_8(37)_{10}(33)_{\mathrm{exp}}\ [0.37\mathrm{ppb}]\,,\\ \alpha^{-1}(\mathrm{a_e}11) &= 137\ 035\ 999\ 1\cdot\cdot\ (09)_8(06)_{10}(33)_{\mathrm{exp}}\ [0.254\mathrm{ppb}]\,. \end{split}
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A few degists in $\alpha^{-1}(a_e11)$ are dotted to avoid the value being quoted.

$(g-2)_{\mu}$: Contributions

- $a_{\mu} = (g-2)_{\mu}/2 = a_{\mu}^{\text{QED}} + a_{\mu}^{\text{EW}} + a_{\mu}^{\text{had}} + a_{\mu}^{\text{New Physics?}}$
- QED: 4-loop predictions consolidated, 5-loop calculations ongoing, big surprises very unprobable for a_{μ} , error formidably small: $a_{\mu}^{\rm QED} = 116584718.08(15) \cdot 10^{-11} \checkmark$ Kinoshita et al.
- EW: consistent 2-loop predictions, accuracy fully sufficient: $a_{\mu}^{\rm EW} = (154 \pm 2) \cdot 10^{-11} \checkmark$ Czarnecki et al., Knecht et al.
- Hadronic contributions: uncertainties completely dominate $\Delta a_{\mu}^{\rm SM}$ imes

$$a_{\mu}^{\mathrm{had}} = a_{\mu}^{\mathrm{had,VP\ LO}} + a_{\mu}^{\mathrm{had,VP\ NLO}} + a_{\mu}^{\mathrm{had,Light-by-Light}}$$

- \blacktriangleright Hadronic contributions from low γ virtualities not calculable with perturbative QCD
- Lattice simulations difficult: accuracy not (yet?!) competetive \rightarrow K. Jansen

$(g-2)_{\mu}$: Contributions

► Hadronic Vacuum Polarisation from exp. $\sigma(e^+e^- \to \gamma^* \to hadrons(+\gamma))$ data [or from $\tau \to \nu_\tau + hadrons$ spectral functions; isospin breaking... \to talk of M. Benayoun] Use of dispersion integral (based on analyticity and unitarity):

$$a_{\mu}^{\text{had,VP LO}} = \frac{1}{4\pi^3} \int_{m_{\pi}^2}^{\infty} ds \, \sigma_{\text{had}}^0(s) K(s) , \quad \text{with } K(s) = \frac{m_{\mu}^2}{3s} \cdot (0.4 \dots 1)$$

- \rightarrow Kernel $K \leadsto$ weighting towards smallest energies. $\sigma_{\rm had}^0$ the undressed cross section
- ightarrow Similar approach with different kernel functions for NLO VP contributions $a_{\mu}^{{
 m had,VP~NLO}}$
- ► Hadronic Light-by-Light:

 \rightarrow see also talk of A. Radzhabov

— No dispersion relation. First Principles calculations from lattice QCD are underway...

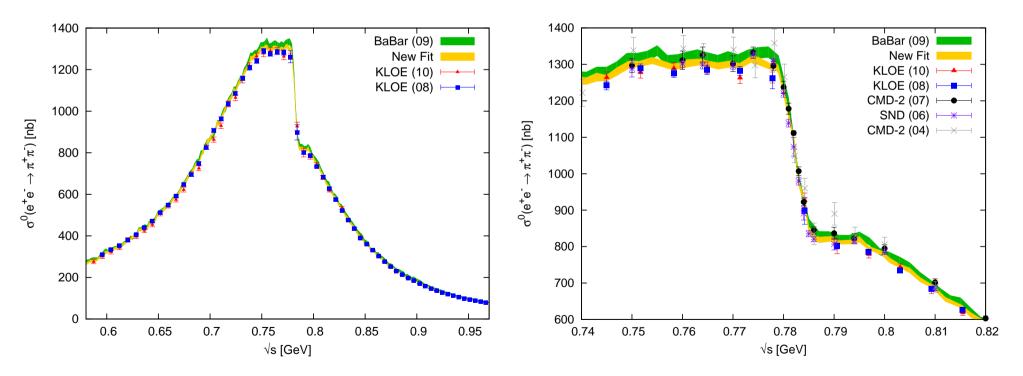
 \rightarrow talk by K. Jansen

Also first results based on Dyson-Schwinger eqs. by C. Fischer et al.

- 'Consensus' of different recent model calculations. HLMNT numbers below use compilation from J. Prades, E. de Rafael, A. Vainshtein: $a_{\mu}^{\rm L-by-L}=(10.5\pm2.6)\cdot10^{-10}$
- Compatible result from F. Jegerlehner, A. Nyffeler: $a_{\mu}^{\rm L-by-L}=(11.6\pm4.0)\cdot10^{-10}$

Recent developments in $(g-2)_{\mu}$; Hadronic VP contributions

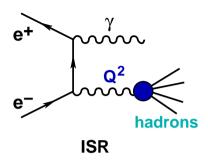
- For low energy $\sigma_{\rm had}^0(s)$, need to sum ~ 25 exclusive channels $[2\pi, 3\pi, KK, 4\pi, \ldots]$
- $\sqrt{s} \sim 1.4 2$ GeV: sum exclusive channels and/or use old inclusive data
- above ~ 2 GeV: inclusive data or use of perturbative QCD [+narrow resonances]
- ▶ The most important 2π channel (> 70%) HLMNT '11 uses 879 data points



Overall, the data combination incl. 'Direct Scan' and 'Radiative Return' looks fine, but...

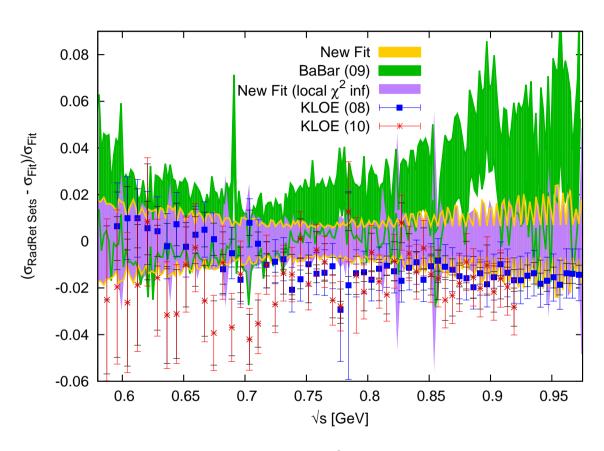
Radiative Return $\pi\pi(\gamma)$ data [KLOE 08/10 and BaBar 09] compared to combination of all

→ talks of P. Lukin, A. Hafner



- \rightarrow Radiative Return (at fixed e^+e^- energy) a powerful method, complementary to direct energy scan
- → Differences in shape and BaBar high at medium and higher energies
- → limited gain in accuracy due to 'tension'; pull-up (mainly from BaBar)

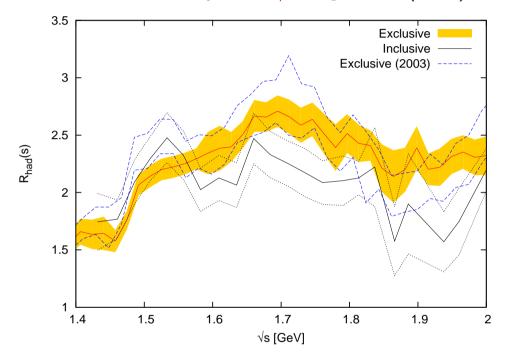
Normalised difference of cross sections [HLMNT '11]



- Comb. of all data on same footing, before integration (purple band): still good $\chi^2_{\rm min}/{\rm d.o.f.} \sim 1.5$ of fit]
- $\bullet \ a_{\mu}^{2\pi}(0.32-2\,\mathrm{GeV}) = (\mathbf{504.2\pm3.0})\cdot 10^{-10}, \ a_{\mu}^{2\pi, \text{w/outRad. Ret.}} = (\mathbf{498.7\pm3.3})\cdot 10^{-10}.$
- Clarification/improvement with more, possibly even more precise data (from both scan and ISR)!?!

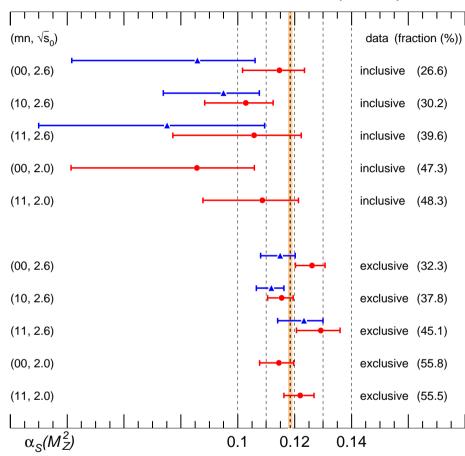
Data compilation $R_{\rm had} = \sigma_{\rm had}/[4\pi\alpha^2/(3s)]$:

blue: old excl. analysis, red/orange: new (2011)



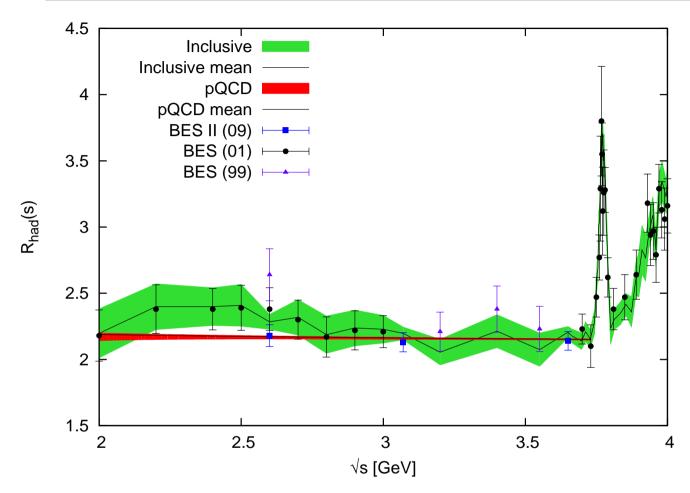
sum over exclusive channels improved

Sum-rules 'determining' α_s (2011):



- shape similar, but normalisation (still) different from old inclusive data
- finite energy QCD sum rules \rightsquigarrow use of exclusive data now preferred over inlcusive
- ullet Still rely on isospin relations for missing channels [sizeable error from $K\bar{K}\pi\pi$]
- For HLMNT '11: Use of more precise sum over exclusive (\hookrightarrow shift up by $\sim +2.7 \cdot 10^{-10}$ for a_{μ})

Perturbative QCD vs. inclusive data above 2 GeV (below the charm threshold)



- Latest BES data (blue markers) in perfect agreement with perturbative QCD; data slightly higher than pQCD for $\sqrt{s} > 2.6$ GeV
- \bullet HLMNT use pQCD for $2.6 < \sqrt{s} < 3.7$ GeV and with (larger) BES errors
 - would have small shift downwards ($\sim -1.4 \cdot 10^{-10}$ for a_{μ}) if used from 2 GeV
 - Davier et al. use pQCD from 1.8 GeV

$a_{\mu}^{ m SM}$: The full picture

[HLMNT '11: J. Phys. G 38 (2011) 085003]

Source	contr. to $a_{\mu} \cdot 10^{11}$	remarks
QED	$116\ 584\ 718.08 \pm 0.15$	up to 5-loop (Kinoshita+Nio, Passera)
EW	154 ± 2	2-loop, Czarnecki+Marciano+Vainshtein
		$(agrees\ very\ well\ with\ Knecht+Peris+Perrottet+deRafael)$
LO hadr.	6923 ± 42	Davier <i>et al.</i> '10 (e^+e^-)
	6908 ± 47	F. Jegerlehner $+$ R. Szafron '11 $\left(e^{+}e^{-} ight)$
	$6894\pm42\pm18$	${\sf Hagiwara+Martin+Nomura+T~'06}$
	$6949\pm37\pm21$	HLMNT '11, as discussed. Combined error is 43
NLO hadr.	$-98.4 \pm 0.6 \pm 0.4$	HLMNT, in agreem. with Krause '97, Alemany $+D+H$ '98
L-by-L	105 ± 26	▶ Prades+deRafael+Vainshtein
agrees with	$<159\ (95\%\ {\rm CL})$	upper bound from Erler+Toledo Sánchez from PHD
< Nov. 2001:	(-85 ± 25)	the 'famous' sign error, $2.6\sigma \rightarrow 1.6\sigma$
$\sum_{i=1}^{n}$	116591828 ± 49	HLMNT '11

The theory prediction of g-2 is now slightly more precise than the BNL measurement

Showdown: SM vs BNL measurement

$$a_{\mu}^{\mathrm{EXP}} = 116\ 592\ 089(63) \cdot 10^{-11}\ (0.5 \mathrm{ppm})$$



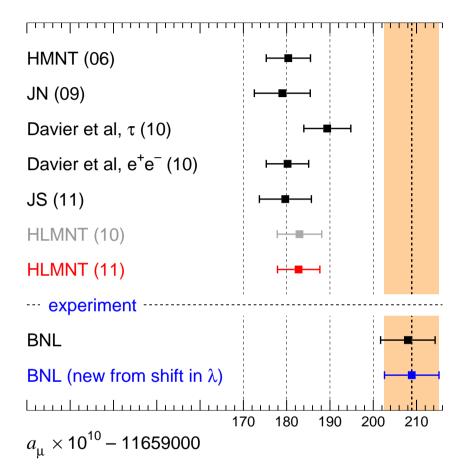
Covered BNL storage ring (Pic. from the g-2 Collab.)

Various choices w.r.t. data, way to compile, τ (?!), L-by-L.

But: ALWAYS had and have found $a_{\mu}^{\mathrm{SM}} < a_{\mu}^{\mathrm{EXP}}$

$$a_{\mu}^{\text{EXP}} - a_{\mu}^{\text{SM}} = (26.1 \pm 8.0) \cdot 10^{-10}$$

 \hookrightarrow 3.3 σ discrepancy [using HLMNT '11]



- Davier et al. '10: 3.6 σ [e^+e^-]
- Jegerlehner+Nyffeler '09: $3.2\,\sigma$
- J+Szafron '11: 3.3 σ [incl. τ w. $\gamma \rho$ mix.]
- HLMNT '09: was 4.0σ [w/out BaBar 09 2π]
- \rightarrow talk by M. Benayoun:

> 4 σ incl. τ w. 'HLS', no ISR 2π data

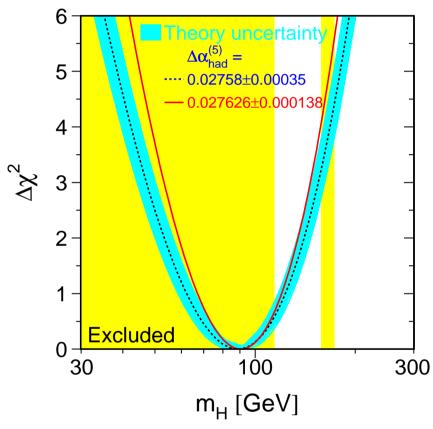
The running QED coupling $\alpha(M_Z^2)$

... and the Higgs mass



- Vacuum polarisation leads to the 'running' of $\alpha \ \text{from} \ \alpha(q^2=0) = 1/137.035999084(51)$ to $\alpha(q^2=M_Z^2) \sim 1/129$
- $\alpha(q^2) = \alpha / (1 \Delta \alpha_{\text{lep}}(q^2) \Delta \alpha_{\text{had}}(q^2))$
- Again use of a dispersion relation: $\Delta\alpha_{\rm had}^{(5)}(q^2) = -\frac{\alpha q^2}{3\pi} P \int_{s_{\rm th}}^{\infty} \frac{R_{\rm had}(s) \, ds}{s(s-q^2)}$
- Hadronic uncertainties \leadsto α the least well known EW param. of $\{G_{\mu}, M_Z, \alpha(M_Z^2)\}$!
- ullet We find (HLMNT '11): $\Deltalpha_{
 m had}^{(5)}(M_Z^2)=0.027626\pm0.000138$ i.e. $lpha(M_Z^2)^{-1}=128.944\pm0.019$
- ullet HLMNT-routine for $lpha(q^2)$ & $R_{
 m had}^{
 m data}$ available

Fit of the SM Higgs mass: LEP EWWG



Fit and Fig. thanks to M. Grünewald

$$\hookrightarrow M_H = 91^{+30}_{-23}\,\mathrm{GeV}$$
 [$m_t = (173.3 \pm 1.1)\,\mathrm{GeV}$]

Outlook

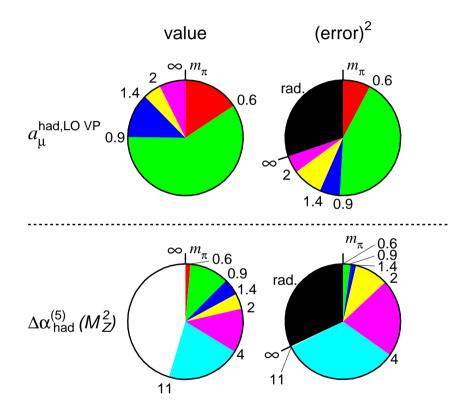
Further improvements

Hadronic VP still the biggest error in a_{μ}^{SM} , soon L-by-L...

Prospects for further squeezing errors:

- More Rad. Ret. in progress at KLOE
- Great opportunity for KLOE-2, BELLE, Super $\tau-c$, in a few years SUPER-Bs, also strong case for DAFNE-HE
- Big improvement envisaged with CMD-3 and SND at VEPP2000
- Higher energies: BES-III at BEPCII in Beijing is on; KEDR at VEPP-4M

Pie diagrams for contr. to a_{μ} and $lpha(M_Z)$ and their errors 2



- ▶ New g-2 experiments at Fermilab and J-PARC.
- → talks by G. Venanzoni, N. Saito
- \blacktriangleright Will a_{μ}^{SM} match the planned accuracy? \leadsto L-by-L may become the limiting factor!

Conclusions

- $(g-2)_{\mu}$ strongly tests all sectors of the SM and constrains possible physics beyond.
- SM prediction consolidated in all sectors: Loops for QED + EW, many exp. data for $R_{\rm had}$ plus TH (incl. Rad. Ret.) for hadronic VP, low energy modelling for L-by-L.
- With the same data compilations as for g-2, also the hadronic contributions to $\Delta\alpha(q^2)$ have been determined; in turn $\alpha(M_Z^2)$ has been improved considerably.
 - \hookrightarrow EW precision fits: $M_H=91^{+30}_{-23}~{
 m GeV}$
- Interaction of TH + MC + EXP most important to achieve even higher precision.
 - → join the WG Radio Montecar Low meeting tomorrow
- low energy $R_{\rm had}$ is also a place to measure α_s at a low scale.
- ightharpoonup Discrepancy betw. SM pred. of g-2 and BNL measurement persists at well $>3\,\sigma$.
- ▶ More to come from all sides. Clear and strong case for continued *and* new experiments!

The coming years will be exciting, and not only for the LHC